
Simple Hyperelastic BVPs

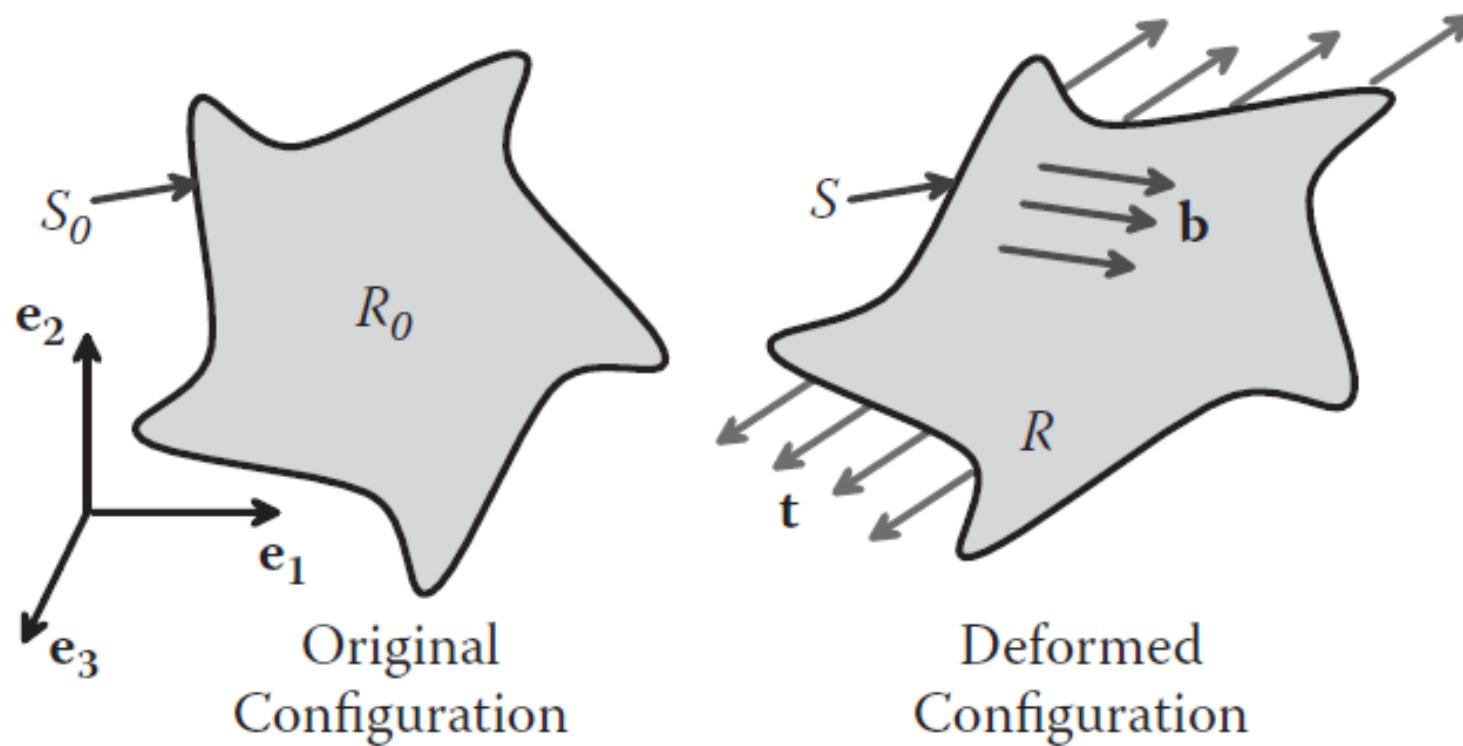
mi@seu.edu.cn

Outline

- Theory of hyperelasticity (超弹理论回顾)
- Incompressible spherically symmetric solids (不可压缩中心对称体)
- Pressurized hollow sphere (压力球腔)
 - Governing equations (控制方程)
 - Boundary condition (边界条件)
 - Displacement vs. Pressure (位移与压力函数关系)
 - Radial stress distribution (径向应力分布)
 - Hoop stress distribution (箍筋应力/周向应力)

Summary of the Theory of Hyperelasticity

- The solid is stress free in its undeformed configuration.
- Temperature changes during deformation are neglected.
- The solid is incompressible.



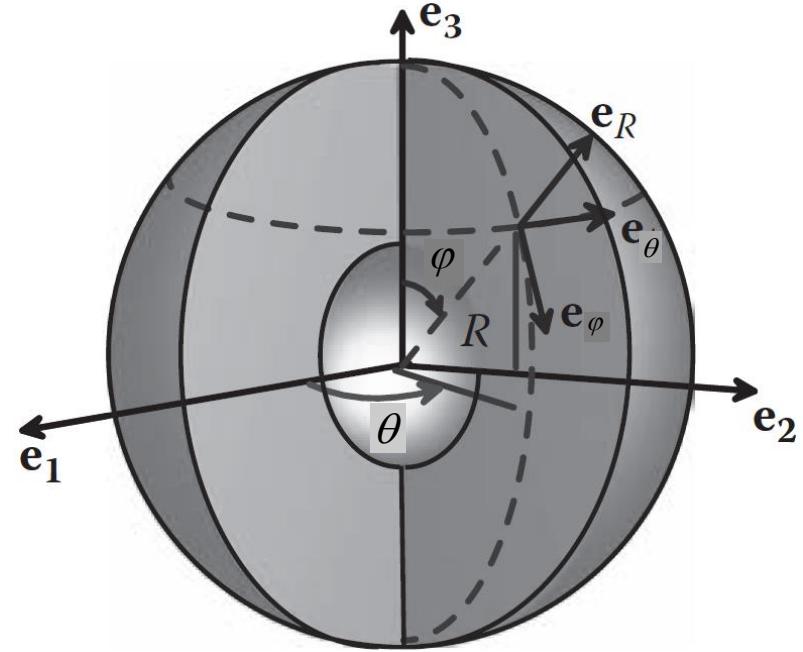
Summary of the Theory of Hyperelasticity

- Strain-displacement relations: $B_{ij} = F_{ik}F_{jk}$, $F_{ij} = \delta_{ij} + u_{i,j}$
- Incompressibility: $J = \det[\mathbf{F}] = 1$
- Stress-strain relation $I_1 = B_{kk}, I_2 = \frac{1}{2}(I_1^2 - B_{ik}B_{ki}),$
 $\sigma_{ij} = 2 \left\{ \left(\frac{\partial U}{\partial I_1} + I_1 \frac{\partial U}{\partial I_2} \right) B_{ij} - \left(I_1 \frac{\partial U}{\partial I_1} + 2I_2 \frac{\partial U}{\partial I_2} \right) \frac{\delta_{ij}}{3} - B_{im}B_{jm} \frac{\partial U}{\partial I_2} \right\} + p\delta_{ij}.$
- Equilibrium equations: $\frac{\partial \sigma_{ij}}{\partial y_i} + F_j = 0.$
- Traction BCs on S_t : $\sigma_{ij}n_i = t_j$
- Displacement BCs on S_u : $u_i = \bar{u}_i$

Incompressible Spherically Symmetric Solids

- Coordinates in **undeformed** configuration $\{R, \Phi, \Theta\}$
- Coordinates in **deformed** configuration $\{r, \varphi, \theta\}$
- Points only move radially, due to spherical symmetry

$$r = f(R) \quad \theta = \Theta \quad \varphi = \Phi.$$



- Position vector in the undeformed solid: $\mathbf{x} = R\mathbf{e}_r$
- Position vector in the deformed solid: $\mathbf{y} = r\mathbf{e}_r = f(R)\mathbf{e}_r$
- Displacement vector: $\mathbf{u} = \mathbf{y} - \mathbf{x} = r\mathbf{e}_r - R\mathbf{e}_r = (f(R) - R)\mathbf{e}_r$

Incompressible Spherically Symmetric Solids

- Cauchy stress: $\boldsymbol{\sigma} = \sigma_{rr}[r]\mathbf{e}_r\mathbf{e}_r + \sigma_{\varphi\varphi}[r]\mathbf{e}_\varphi\mathbf{e}_\varphi + \sigma_{\theta\theta}[r]\mathbf{e}_\theta\mathbf{e}_\theta, \quad \sigma_{\varphi\varphi} = \sigma_{\theta\theta}.$

- Deformation gradient:

$$\mathbf{F} = F_{rr}[r, R]\mathbf{e}_r\mathbf{e}_R + F_{\varphi\varphi}[r, R]\mathbf{e}_\varphi\mathbf{e}_\Phi + F_{\theta\theta}[r, R]\mathbf{e}_\theta\mathbf{e}_\Theta, \quad F_{\varphi\varphi} = F_{\theta\theta}.$$

- Left C-G deformation tensor:

$$\mathbf{B} = B_{rr}[r]\mathbf{e}_r\mathbf{e}_r + B_{\varphi\varphi}[r]\mathbf{e}_\varphi\mathbf{e}_\varphi + B_{\theta\theta}[r]\mathbf{e}_\theta\mathbf{e}_\theta, \quad B_{\varphi\varphi} = B_{\theta\theta}.$$

- Strain-displacement relations:

$$F_{rr} = \frac{dr}{dR}, \quad F_{\varphi\varphi} = F_{\theta\theta} = \frac{r}{R}, \quad \Rightarrow B_{rr} = \left(\frac{dr}{dR} \right)^2, \quad B_{\varphi\varphi} = B_{\theta\theta} = \left(\frac{r}{R} \right)^2$$

- Incompressibility:

$$J = \det[\mathbf{F}] = \frac{dr}{dR} \left(\frac{r}{R} \right)^2 = 1.$$

Incompressible Spherically Symmetric Solids

- Stress-strain relation

$$\sigma_{rr} = 2 \left\{ \left(\frac{\partial U}{\partial I_1} + I_1 \frac{\partial U}{\partial I_2} \right) B_{rr} - \frac{1}{3} \left(I_1 \frac{\partial U}{\partial I_1} + 2I_2 \frac{\partial U}{\partial I_2} \right) - B_{rr}^2 \frac{\partial U}{\partial I_2} \right\} + p,$$

$$\sigma_{\varphi\varphi} = \sigma_{\theta\theta} = 2 \left\{ \left(\frac{\partial U}{\partial I_1} + I_1 \frac{\partial U}{\partial I_2} \right) B_{\varphi\varphi} - \frac{1}{3} \left(I_1 \frac{\partial U}{\partial I_1} + 2I_2 \frac{\partial U}{\partial I_2} \right) - B_{\varphi\varphi}^2 \frac{\partial U}{\partial I_2} \right\} + p.$$

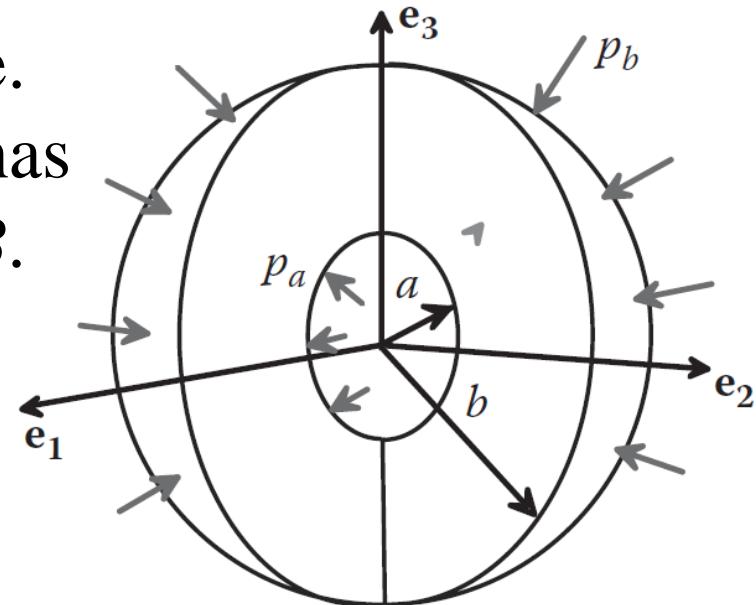
- Equilibrium equations: $\frac{d\sigma_{rr}}{dr} + \frac{2}{r} (\sigma_{rr} - \sigma_{\varphi\varphi}) + F_r = 0.$
- Traction BCs: $\sigma_{rr}[a] = \sigma_a, \sigma_{rr}[b] = \sigma_b.$
- Displacement BCs: $u_r[a] = u_a, u_r[b] = u_b.$

Pressurized Hollow Sphere: Incompressibility

- No body forces act on the sphere.
- Before deformation, the sphere has inner radius A and outer radius B .
- The solid is made from an incompressible Mooney-Rivlin solid.
- Solution:
- Incompressibility implies:

$$V = V_0 \quad \Rightarrow \int_a^r 4\pi r^2 dr = \int_A^R 4\pi R^2 dR \quad \Rightarrow r^3 - a^3 = R^3 - A^3$$

$$\Rightarrow r = (R^3 - A^3 + a^3)^{\frac{1}{3}}, \quad R = (r^3 - a^3 + A^3)^{\frac{1}{3}}$$



Pressurized Hollow Sphere: Governing Equations

- Deformation gradient and left C-G strain:

$$F_{rr} = \frac{dr}{dR} = \left(\frac{r}{R} \right)^2, F_{\varphi\varphi} = F_{\theta\theta} = \frac{r}{R}, \Rightarrow B_{rr} = \left(\frac{r}{R} \right)^4, B_{\varphi\varphi} = B_{\theta\theta} = \left(\frac{r}{R} \right)^2$$

- Stress-strain relations:

$$I_1 = B_{kk} = B_{rr} + 2B_{\varphi\varphi}, \quad I_2 = \frac{1}{2} \left(I_1^2 - B_{ik} B_{ki} \right) = \frac{1}{2} \left\{ \left(B_{rr} + 2B_{\varphi\varphi} \right)^2 - \left(B_{rr}^2 + 2B_{\varphi\varphi}^2 \right) \right\} = \left(2B_{rr} + B_{\varphi\varphi} \right) B_{\varphi\varphi}$$

$$U = \frac{\mu_1}{2} (I_1 - 3) + \frac{\mu_2}{2} (I_2 - 3) \quad \Rightarrow \frac{\partial U}{\partial I_1} = \frac{\mu_1}{2}, \quad \frac{\partial U}{\partial I_2} = \frac{\mu_2}{2}$$

$$\Rightarrow \boxed{\sigma_{rr} = \frac{2}{3} (\mu_1 + \mu_2 B_{\varphi\varphi}) (B_{rr} - B_{\varphi\varphi}) + p, \quad \sigma_{\varphi\varphi} = \sigma_{\theta\theta} = -\frac{1}{3} (\mu_1 + \mu_2 B_{\varphi\varphi}) (B_{rr} - B_{\varphi\varphi}) + p.}$$

- Equilibrium equation:

$$\frac{d\sigma_{rr}}{dr} + \frac{2}{r} (\sigma_{rr} - \sigma_{\varphi\varphi}) + F_r = 0 \quad \Rightarrow \sigma_{rr} = \mu_1 \left(\frac{2R}{r} + \frac{R^4}{2r^4} \right) - \mu_2 \left(\frac{2r}{R} - \frac{R^2}{r^2} \right) + C.$$

Pressurized Hollow Sphere: BCs

- Boundary conditions

$$\begin{cases} r = a; R = A : \sigma_{rr} = -p_a \\ r = b; R = B : \sigma_{rr} = -p_b \end{cases} \Rightarrow \begin{cases} -p_a = \mu_1 \left(\frac{2A}{a} + \frac{A^4}{2a^4} \right) - \mu_2 \left(\frac{2\alpha}{A} - \frac{A^2}{a^2} \right) + C \equiv \mu_1 \left(\frac{2}{\alpha} + \frac{1}{2\alpha^4} \right) - \mu_2 \left(2\alpha - \frac{1}{\alpha^2} \right) + C \\ -p_b = \mu_1 \left(\frac{2B}{b} + \frac{B^4}{2b^4} \right) - \mu_2 \left(\frac{2\beta}{B} - \frac{B^2}{b^2} \right) + C \equiv \mu_1 \left(\frac{2}{\beta} + \frac{1}{2\beta^4} \right) - \mu_2 \left(2\beta - \frac{1}{\beta^2} \right) + C \end{cases}$$

- Adding the two equations gives the expression for C .

$$C = -\frac{\mu_1}{2} \left(\frac{2}{\alpha} + \frac{1}{2\alpha^4} + \frac{2}{\beta} + \frac{1}{2\beta^4} \right) + \frac{\mu_2}{2} \left(2\alpha - \frac{1}{\alpha^2} + 2\beta - \frac{1}{\beta^2} \right) - \frac{p_a + p_b}{2}$$

- Although unnecessary, the Pressure p in the stress-strain relations can be determined by comparing expressions of the radial stress.

Pressurized Hollow Sphere: Displacement vs. Pressure

- Subtracting the first BC equation from the second results in the expression that relates α and β to the pressures.

$$\frac{p_a - p_b}{\mu_1} = 2\left(\frac{1}{\beta} - \frac{1}{\alpha}\right) + \frac{1}{2}\left(\frac{1}{\beta^4} - \frac{1}{\alpha^4}\right) - \frac{2\mu_2}{\mu_1}(\beta - \alpha) + \frac{\mu_2}{\mu_1}\left(\frac{1}{\beta^2} - \frac{1}{\alpha^2}\right)$$

- Incompressibility yields

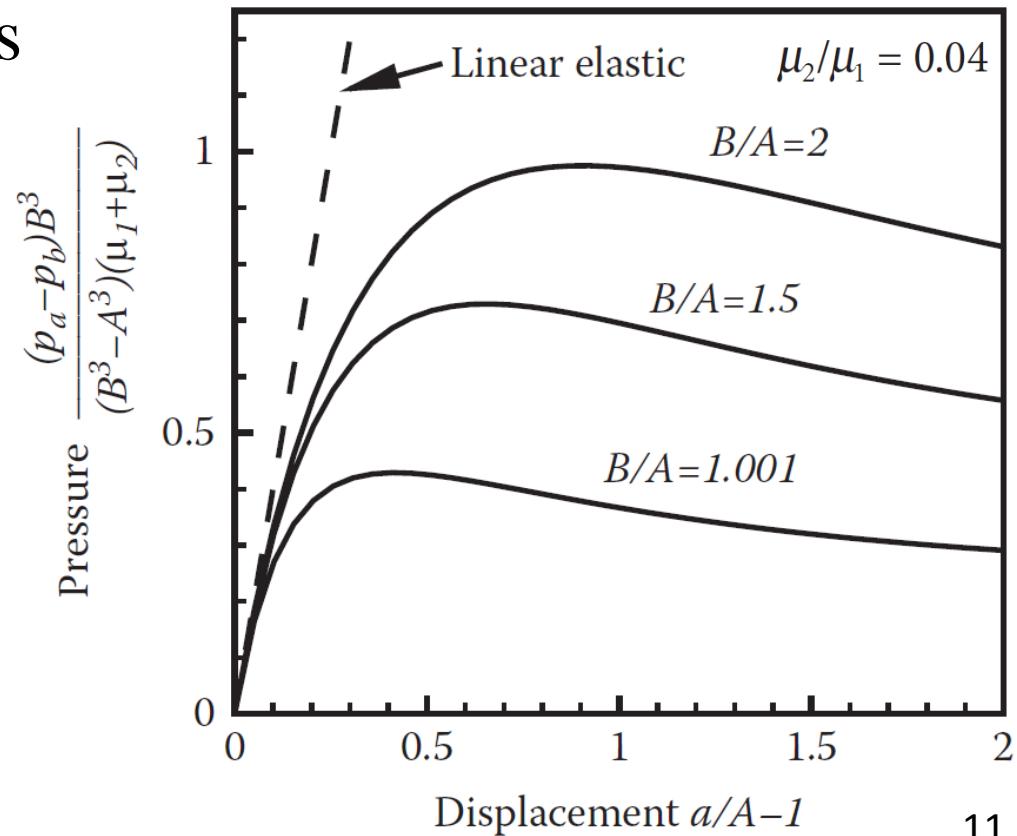
$$r^3 - a^3 = R^3 - A^3$$

$$\Rightarrow A^3 - a^3 = R^3 - r^3 = B^3 - b^3$$

$$\Rightarrow \frac{B^3}{A^3} = \frac{B^3}{A^3} \frac{A^3 - a^3}{B^3 - b^3} = \frac{1 - \alpha^3}{1 - \beta^3}$$

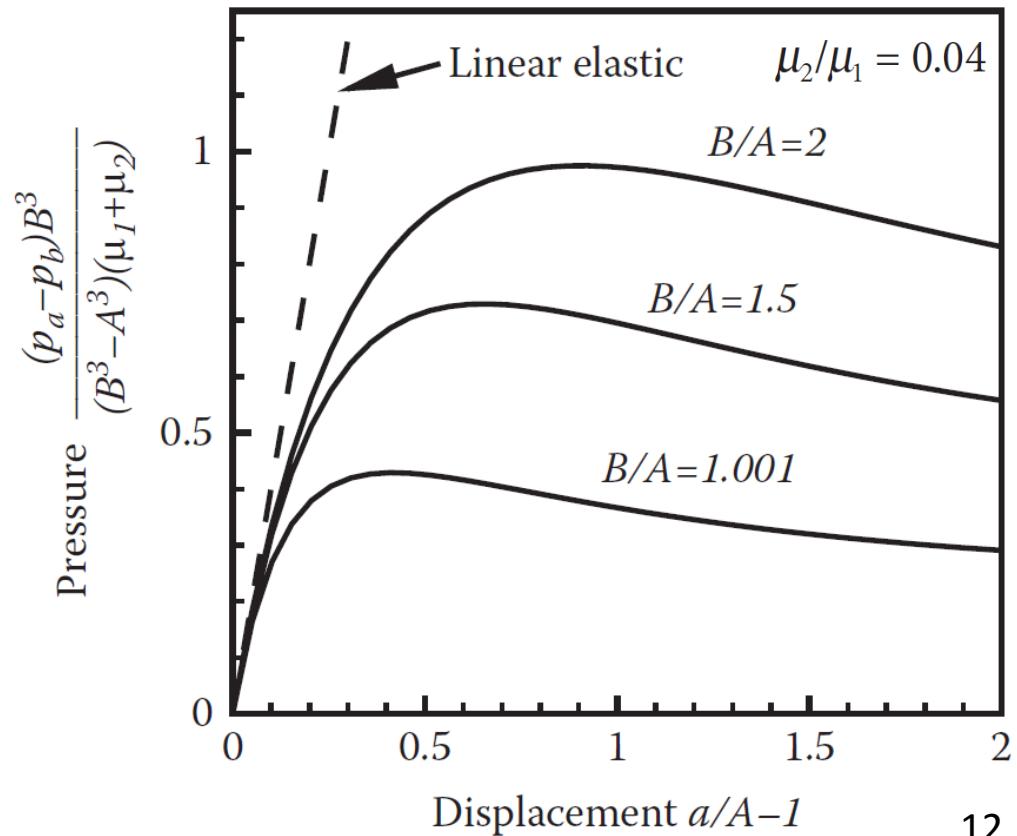
- The two equations can be solved for α and β .
- For graphing purpose, it is better to follow:

$$\alpha \rightarrow \beta \rightarrow p_a - p_b.$$



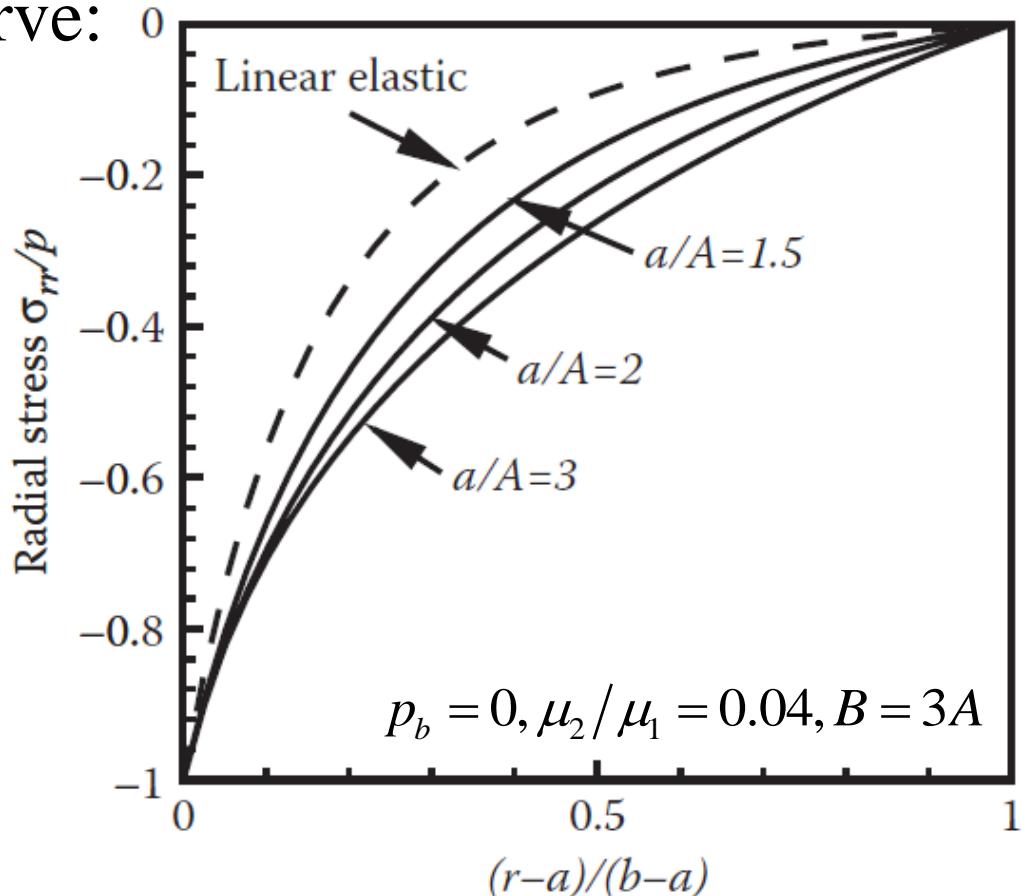
Pressurized Hollow Sphere: Displacement vs. Pressure

- For the linear elastic curve: $E = 3(\mu_1 + \mu_2)$, $\nu = 0.5$.
- The linear elastic solution is accurate for $u/A < 0.05$.
- The relationship between pressure and displacement is nonlinear in the large deformation regimen.
- As the internal radius of the sphere increases, the pressure reaches a maximum and thereafter decreases because the wall thickness of the shell decreases as the sphere expands.



Pressurized Hollow Sphere: Radial Stress

- Radial stress distribution: $\sigma_{rr} = \mu_1 \left(\frac{2R}{r} + \frac{R^4}{2r^4} \right) - \mu_2 \left(\frac{2r}{R} - \frac{R^2}{r^2} \right) + C$
- For the linear elastic curve: $E = 3(\mu_1 + \mu_2)$, $\nu = 0.5$.
- The radial stress remains close to the linear elastic solution even in the large deformation regimen.



Pressurized Hollow Sphere: Hoop Stress

- Hoop stress

$$\sigma_{rr} = \frac{2}{3}(\mu_1 + \mu_2 B_{\varphi\varphi})(B_{rr} - B_{\varphi\varphi}) + p, \quad \sigma_{\varphi\varphi} = \sigma_{\theta\theta} = -\frac{1}{3}(\mu_1 + \mu_2 B_{\varphi\varphi})(B_{rr} - B_{\varphi\varphi}) + p.$$

$$\Rightarrow \sigma_{\varphi\varphi} = \sigma_{rr} - (\mu_1 + \mu_2 B_{\varphi\varphi})(B_{rr} - B_{\varphi\varphi}) \\ = \mu_1 \left(\frac{2R}{r} - \frac{R^4}{2r^4} + \frac{r^2}{R^2} \right) - \mu_2 \left(\frac{2r}{R} - \frac{r^4}{R^4} \right) + C$$

- For the linear elastic curve:

$$E = 3(\mu_1 + \mu_2), \nu = 0.5.$$

- The hoop stress distribution is significantly altered as the deformation increases

